# $L^{p}$-SOLUTIONS OF FRACTIONAL DIFFERENTIAL EQUATIONS 

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#### Abstract

We study fractional differential equations of Caputo type ${ }^{c} D^{q} x(t)=u(t, x(t)), 0<q<1$, of both linear and nonlinear type. That equation is inverted as an integral equation with kernel $C(t-s):=(1 / \Gamma(q))(t-s)^{q-1}$. We then transform the integral equation into one with kernel $R(t-s)$ so that $0<R(t) \leq C(t)$ and $\int_{0}^{\infty} R(s) d s=1$. A variety of techniques are introduced by which we are able to show that solutions are in $L^{p}[0, \infty)$ for appropriate $p \geq 1$.


## 1. Introduction

The object of this paper is to present basic techniques for showing that the solutions of common fractional differential equations are in $L^{p}[0, \infty)$ for some positive integer $p$. Here is a loose description of how we will proceed. Invert the fractional differential equation of Caputo type

$$
{ }^{c} D^{q} x(t)=u(t, x(t)), \quad 0<q<1, \quad x(0) \in \Re
$$

as the standard integral equation

$$
x(t)=x(0)+\frac{1}{\Gamma(q)} \int_{0}^{t}(t-s)^{q-1} u(s, x(s)) d s
$$

where $\Gamma$ is the gamma function. Because of the large and singular kernel it can be very difficult to study. Moreover, the $x(0)$ is a constant source of difficulty. There are at least two reasonable ways to avoid those problems. We use both.

Stability theory assumes that we have an object in equilibrium, $x(0)=0$, and we disturb the equilibrium with $x(0) \neq 0$, causing no further disturbance. We then study the subsequent position of the object.

An equally important study assumes that $x(0)=0$, but we constantly perturb the object by an external force, say $f(t)$, and study the resulting position of the object. Three very common forms of $f(t)$ are:
(i) $f(t) \rightarrow 0$ as $t \rightarrow \infty$.
(ii) $f \in L^{p}[0, \infty)$.
(iii) $f$ is periodic.

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Our choice here is (ii) and we study both $x(0)=0$ and $x(0) \neq 0$. Moreover, the reader may consult Podlubny [7, p. 224] to see that the assumption of $x(0)=0$ is of general use and not simply contrived to meet our requirements. For that case we study

$$
{ }^{c} D^{q} x(t)=u(t, x(t))+f(t), \quad 0<q<1, \quad f \in L^{p}[0, \infty), \quad x(0)=0 .
$$

The inversion is then

$$
x(t)=\frac{1}{\Gamma(q)} \int_{0}^{t}(t-s)^{q-1}[u(s, x(s))+f(s)] d s
$$

and we are thwarted because $f \in L^{1}[0, \infty)$ yields $\int_{0}^{t}(t-s)^{q-1} f(s) d s \rightarrow 0$ as $t \rightarrow \infty$, but generally not in $L^{p}[0, \infty)$.

Fortunately, the kernel is a completely monotone function and that means that the resolvent kernel, $R(t)$, is positive, $\int_{0}^{\infty} R(t) d t=1$, and completely monotone. But to really utilize this wonderful resolvent we need one assumption: $u(t, x)$ must contain a linear term " $-x$ ", Much of the point here is that we can successfully render this true by a variety of techniques, not the least of which is the elementary device of writing $u(t, x)=-x+x+u(t, x)$.

With this assumption in hand we decompose our equation into a pair:

$$
z(t)=x(0)-\int_{0}^{t} R(t-s) x(0) d s
$$

and

$$
x(t)=z(t)+\int_{0}^{t} R(t-s)[x(s)+u(s, x(s))+f(s)] d s .
$$

Here, $z(t) \rightarrow 0$ as $t \rightarrow \infty$; or $z(t)=0$ if $x(0)=0$. Moreover, $\int_{0}^{t} R(t-$ s) $f(s) d s \in L^{1}[0, \infty)$ if $f \in L^{1}[0, \infty)$. This gives us a real chance of proving $x \in L^{p}[0, \infty)$, a chance that was essentially zero in the original form.

In the next pages we offer a set of examples showing how this can be done. We will introduce a positive constant, $J$, which preserves the complete monotonicity of the kernel and makes a number of things possible. Not the least of these is that mappings which were only Lipschitz become contractions.

## 2. BASIC THEORY

Here are the details of the decomposition described in Section 1. The inversion of the Caputo equation into an integral equation when $u(t, x)$ is continuous is verified in ([4, p. 54], [3, pp. 78, 86, 103]). This equation has a unique solution as seen in [1] or [8]). We begin with

$$
\begin{equation*}
{ }^{c} D^{q} x(t)=u(t, x), \quad 0<q<1, \quad x(0) \in \Re \tag{1}
\end{equation*}
$$

where $u:[0, \infty) \times \Re \rightarrow \Re$ is continuous. At this point a positive constant $J$ is introduced and we invert (1) as

$$
x(t)=x(0)+\frac{J}{\Gamma(q)} \int_{0}^{t}(t-s)^{q-1} \frac{u(s, x(s))}{J} d s
$$

In our subsequent work, the function " $-x$ " is introduced in a wellmotivated way. But here we simply add and subtract to obtain
(2) $x(t)=x(0)+\frac{J}{\Gamma(q)} \int_{0}^{t}(t-s)^{q-1}\left[-x(s)+x(s)+\frac{u(s, x(s))}{J}\right] d s$

Denote the kernel by

$$
\begin{equation*}
C(t)=\frac{J}{\Gamma(q)} t^{q-1} \tag{3}
\end{equation*}
$$

so that for any $T>0$ we have the critical property that

$$
\int_{0}^{T}|C(u)| d u<\infty
$$

Following Miller [6, pp. 193-22] we note that $C(t)$ is completely monotone on $(0, \infty)$ in the sense that $(-1)^{k} C^{(k)}(t) \geq 0$ for $k=0,1,2, \ldots$ and $t \in(0, \infty)$. Moreover $C(t)$ satisfies the conditions of Miller's Theorem 6.2 on p. 212. That theorem states that if the resolvent equation for the completely monotone kernel $C$ is

$$
\begin{equation*}
R(t)=C(t)-\int_{0}^{t} C(t-s) R(s) d s \tag{4}
\end{equation*}
$$

then that resolvent kernel, $R$, satisfies
(5) $0 \leq R(t) \leq C(t)$ for all $t>0$ so that as $t \rightarrow \infty$ then $R(t) \rightarrow 0$
and that

$$
\begin{equation*}
C \notin L^{1}[0, \infty) \quad \Longrightarrow \int_{0}^{\infty} R(s) d s=1 \tag{6}
\end{equation*}
$$

Continuing on to [6, pp. 221-224 (Theorem 7.2)] we see that $R$ is also completely monotone.

Next, under the conditions here, it is shown in Miller [6, pp. 191-207] that (2) can be decomposed into

$$
\begin{equation*}
z(t)=x(0)-\int_{0}^{t} C(t-s) z(s) d s \tag{7}
\end{equation*}
$$

with

$$
z(t)=x(0)-\int_{0}^{t} R(t-s) x(0) d s=x(0)\left[1-\int_{0}^{t} R(s) d s\right]
$$

and, having found $z(t)$, then the solution $x(t)$ of (2) solves

$$
\begin{equation*}
x(t)=z(t)+\int_{0}^{t} R(t-s)\left[x(s)+\frac{u(s, x(s))}{J}\right] d s \tag{8}
\end{equation*}
$$

Notice that $z(t) \rightarrow 0$ as $t \rightarrow \infty$.
The kernel in (2) is not integrable on $[0, \infty)$, but in (8) it is replaced, not only by an integrable kernel, but the value of the integral is one and the new kernel is also completely monotone.

## 3. $f \in L^{1}[0, \infty)$ : Elementary arguments

Our first theorem, together with the preparation for it, can be viewed as an introduction to and a simple and transparent example of several of our subsequent results. It shows the basic selection of the constant $J$ which is one of the most essential parts of the study. The linear equation

$$
\begin{equation*}
{ }^{c} D^{q} x=f(t)-a(t) x(t), \quad 0<q<1, \quad x(0)=0 \tag{9}
\end{equation*}
$$

is inverted as

$$
\begin{equation*}
x(t)=x(0)-\frac{1}{\Gamma(q)} \int_{0}^{t}(t-s)^{q-1}[a(s) x(s)-f(s)] d s \tag{10}
\end{equation*}
$$

where $a, f:[0, \infty) \rightarrow \Re$ are continuous and there are positive numbers $\epsilon$ and $M$ with

$$
\begin{equation*}
0<\epsilon \leq a(t) \leq M \tag{11}
\end{equation*}
$$

We will exchange the kernel in (10) for $R(t-s)$, but first we will reduce $a(t)$ to a function bounded by $\alpha<1$. Define $J=\epsilon+(1 / 2)(M-\epsilon)$. Then there is an $\alpha$ with

$$
\begin{equation*}
J>0, \quad 0<\alpha<1, \quad|a(t)-J|<\alpha J . \tag{12}
\end{equation*}
$$

Note that we may choose $\alpha=(M-\epsilon) /(M+\epsilon)$. In fact, if we write $J=(M+\epsilon) / 2$, then by (11) we have

$$
\epsilon-J \leq a(t)-J \leq M-J
$$

This implies that

$$
-\frac{1}{2}(M-\epsilon) \leq a(t)-J \leq \frac{1}{2}(M-\epsilon)
$$

and so

$$
|a(t)-J| \leq \frac{1}{2}(M-\epsilon)=\frac{M-\epsilon}{M+\epsilon} J=: \alpha J
$$

Since $x(0)=0$, we write (10) as

$$
\begin{aligned}
x(t) & =-\frac{1}{\Gamma(q)} \int_{0}^{t}(t-s)^{q-1}[J x(s)+(a(s)-J) x(s)-f(s)] d s \\
& =-\frac{1}{\Gamma(q)} \int_{0}^{t} J(t-s)^{q-1}\left[x(s)+\frac{(a(s)-J)}{J} x(s)-\frac{f(s)}{J}\right] d s \\
& =:-\frac{1}{\Gamma(q)} \int_{0}^{t} J(t-s)^{q-1}\left[x(s)+\beta(s) x(s)-\frac{f(s)}{J}\right] d s
\end{aligned}
$$

where $|\beta(t)| \leq \alpha<1$.

The kernel is still completely monotone so the resolvent has exactly the same properties as the $R$ we discussed before, so we retain the name of $R$. Next, decompose this equation into $z(t)=0$ and

$$
\begin{equation*}
x(t)=F(t)-\int_{0}^{t} R(t-s)[\beta(s) x(s)] d s \tag{13}
\end{equation*}
$$

with

$$
F(t)=\frac{1}{J} \int_{0}^{t} R(t-s) f(s) d s, \quad \beta(s)=\frac{a(s)-J}{J}
$$

Notice that $f \in L^{1}$ implies $F \in L^{1}$.
Condition (11) will be critical in Section 4 for general $p$, but we will see in two steps that it can be changed to $-\infty \leq a(t) \leq \infty$.

Theorem 3.1. If (11) holds and if $f \in L^{1}[0, \infty)$, then the solution of (9) is in $L^{1}[0, \infty)$.

Proof. We have

$$
\begin{aligned}
\int_{0}^{t}|x(s)| d s & \leq \int_{0}^{t}|F(s)| d s+\int_{0}^{t} \int_{0}^{u} R(u-s) \alpha|x(s)| d s d u \\
& =\int_{0}^{t}|F(s)| d s+\int_{0}^{t} \int_{s}^{t} R(u-s) d u \alpha|x(s)| d s \\
& \leq \int_{0}^{t}|F(s)| d s+\int_{0}^{t} \alpha|x(s)| d s
\end{aligned}
$$

so that

$$
(1-\alpha) \int_{0}^{t}|x(s)| d s \leq \int_{0}^{t}|F(s)| d s
$$

as required.
In many problems we find that $\beta(t)$ in the above proof can be arbitrarily large for small values of $t$, but eventually it is dominated by $\alpha<1$. This can happen in two simple ways. First, it may be that $a(t)$ is asymptotically periodic in the sense that $a(t)=b(t)+c(t)$ where $b(t)$ is a positive periodic function, while $c(t) \rightarrow 0$ as $t \rightarrow \infty$. In a more complicated way, it may happen that $b(t)$ is again a positive periodic function, but $c \in L^{1}[0, \infty)$. In this case, $\int_{0}^{t} R(t-s) c(s) d s \rightarrow 0$ as $t \rightarrow \infty$ and some algebraic work must be done to bring us to the desired conclusion. The following lemma allows us to avoid those difficulties.

Lemma 3.2. If $G(t)$ is continuous on $[0, \infty)$, then for $\int_{0}^{t / 2} R(v) d v>$ $1 / 2$ it follows that

$$
(1 / 2) \int_{0}^{t / 2}|G(s)| d s \leq \int_{0}^{t} \int_{0}^{u} R(u-s)|G(s)| d s d u
$$

Proof. We have

$$
\begin{aligned}
(1 / 2) \int_{0}^{t / 2}|G(s)| d s & \leq \int_{0}^{t / 2} \int_{0}^{t / 2} R(v) d v|G(s)| d s \\
& \leq \int_{0}^{t / 2} \int_{0}^{t-s} R(v) d v|G(s)| d s \\
& =\int_{0}^{t / 2} \int_{s}^{t} R(u-s) d u|G(s)| d s \\
& \leq \int_{0}^{t} \int_{s}^{t} R(u-s) d u|G(s)| d s \\
& =\int_{0}^{t} \int_{0}^{u} R(u-s)|G(s)| d s d u
\end{aligned}
$$

In the next result, if $K=0$ then it is essentially Theorem 3.1 and the lemma is not used.

Theorem 3.3. Suppose that there is a number $K>0$ and (11) holds for $t \geq K$ and that $f \in L^{1}[0, \infty)$. Then the solution of (9) is in $L^{1}[0, \infty)$.
Proof. Define $J, \alpha, \beta(t)$ as in (12) and (13). Let $\beta^{*}:[0, \infty) \rightarrow \Re$ be a continuous function with $0 \leq \beta^{*}(t) \leq 1-\alpha$ if $0 \leq t \leq K$ and $\beta^{*}(t)=1-\alpha$ if $t>K$. We have from (13) that

$$
\begin{aligned}
& \int_{0}^{t}|x(s)| d s \leq \int_{0}^{t}|F(s)| d s+\int_{0}^{t} \int_{0}^{u} R(u-s)|\beta(s)||x(s)| d s d u \\
& \leq \int_{0}^{t}|F(s)| d s+\int_{0}^{K} \int_{0}^{u} R(u-s)|\beta(s)||x(s)| d s d u \\
& +\int_{K}^{t} \int_{0}^{u} R(u-s)|x(s)-(1-\alpha) x(s)| d s d u \\
& =: \mu(t)+\int_{K}^{t} \int_{0}^{u} R(u-s)[|x(s)|-(1-\alpha)|x(s)|] d s d u \\
& \leq \mu(t)+\int_{0}^{t} \int_{0}^{u} R(u-s)|x(s)| d s d u-\int_{0}^{t} \int_{0}^{u} R(u-s) \beta^{*}(s)|x(s)| d s d u .
\end{aligned}
$$

Interchange the order of integration in the next-to-last integral and then cancel it against the term on the left-hand-side of the display. This yields

$$
\int_{0}^{t} \int_{0}^{u} R(u-s) \beta^{*}(s)|x(s)| d s d u \leq \mu(t)
$$

By Lemma 3.2 if $t$ is so large that $\int_{0}^{t} R(u) d u>1 / 2$ and $t>2 K$ then

$$
(1 / 2) \int_{K}^{t / 2}(1-\alpha)|x(s)| d s \leq(1 / 2) \int_{0}^{t / 2} \beta^{*}(s)|x(s)| d s \leq \mu(t),
$$

as required.
Using the lemma we can now let $\epsilon=0$ in (11). Here, $M$ will take the place of $J$.

Theorem 3.4. Let

$$
{ }^{c} D^{q} x=-a(t) h(t, x)+f(t), x(0)=0,
$$

with $f \in L^{1}[0, \infty), 0 \leq a(t) \leq M$, and $|h(t, x)| \leq|x|$ for all $(t, x), t \geq 0$. If, in addition, $x h(t, x) \geq 0$, then any solution satisfies

$$
\int_{0}^{t / 2}|(a(s) / M) h(s, x(s))| d s \leq 2 \int_{0}^{t}|\widetilde{F}(s)| d s
$$

where $t$ is large enough that $\int_{0}^{t} R(s) d s>1 / 2$ and

$$
\widetilde{F}(t)=\frac{1}{M} \int_{0}^{t} R(t-s) f(s) d s
$$

Proof. Invert the equation as

$$
\begin{aligned}
x(t) & =-\frac{1}{\Gamma(q)} \int_{0}^{t}(t-s)^{q-1}[a(s) h(s, x(s))-f(s)] d s \\
& =-\frac{M}{\Gamma(q)} \int_{0}^{t}(t-s)^{q-1}[(a(s) / M) h(s, x(s))-(f(s) / M)] d s .
\end{aligned}
$$

Then we add and subtract $x(s)$ in the integrand and decompose in the usual way. As $x(0)=0$ we have $z(t)=0$ and our equation is

$$
x(t)=\int_{0}^{t} R(t-s)\left[x(s)-\frac{a(s)}{M} h(s, x(s))+\frac{f(s)}{M}\right] d s .
$$

Here we again retain the name of $R(t)$. Notice that

$$
\left|x(s)-\frac{a(s)}{M} h(s, x(s))\right|=|x(s)|-\left|\frac{a(s)}{M} h(s, x(s))\right| .
$$

Now, take absolute values of both sides so that we have (for $\widetilde{F}(t)$ defined above)

$$
|x(t)| \leq|\widetilde{F}(t)|+\int_{0}^{t} R(t-s)[|x(s)-|(a(s) / M) h(s, x(s))|] d s
$$

Integrate both sides from 0 to $t$. Interchange the order of integration in the first integral and cancel it with the left-hand-side. Apply the result in the previous proof to obtain

$$
\int_{0}^{t / 2}|(a(s) / M) h(s, x(s))| d s \leq 2 \int_{0}^{t}|\widetilde{F}(s)| d s
$$

as required.
We can relate the next theorem to Theorem 3.1 by noting that $A=\epsilon$ and $B=M ; \frac{\partial G(t, x)}{\partial x}=a(t)$.

Theorem 3.5. Suppose there are positive constants $A$ and $B$ so that for

$$
{ }^{c} D^{q} x=u(t, x)=:-G(t, x)+f(t), \quad x(0) \in \Re
$$

we have $f \in L^{1}[0, \infty), G(t, x) \rightarrow 0$ as $x \rightarrow 0, f, G$ continuous,

$$
A \leq \frac{\partial G(t, x)}{\partial x} \leq B, \quad 0 \leq t<\infty
$$

Then there is a unique solution tending to zero. If, in addition, $x(0)=$ 0 and for $|x|$ sufficiently small there is a $\beta<1$ with

$$
\left|x-\frac{G(t, x)}{J}\right| \leq \beta|x|
$$

for some $J>0$, then that unique solution is in $L^{1}[0, \infty)$.
Proof. The equation is inverted as

$$
\begin{aligned}
x(t) & =x(0)+\frac{J}{\Gamma(q)} \int_{0}^{t}(t-s)^{q-1}\left[-\frac{G(s, x(s))}{J}+\frac{f(s)}{J}\right] d s \\
& =\frac{J}{\Gamma(q)} \int_{0}^{t}(t-s)^{q-1}\left[-x(s)+x(s)-\frac{G(s, x(s))}{J}+\frac{f(s)}{J}\right] d s
\end{aligned}
$$

so that $z(t)=x(0)\left[1-\int_{0}^{t} R(s) d s\right]$ and

$$
x(t)=z(t)+F(t)+\int_{0}^{t} R(t-s)\left[x(s)-\frac{G(s, x(s))}{J}\right] d s
$$

with

$$
F(t)=(1 / J) \int_{0}^{t} R(t-s) f(s) d s
$$

It follows that $F \in L^{1}[0, \infty)$ and $F(t) \rightarrow 0$ as $t \rightarrow \infty$. Now

$$
\frac{\partial}{\partial x}\left[x-\frac{G(t, x)}{J}\right]=1-\frac{\partial G(t, x))}{J \partial x}
$$

and

$$
0<A \leq \frac{\partial G(t, x))}{\partial x} \leq B
$$

so take $J$ so large that

$$
0<1-(B / J)=: \alpha_{1}<1
$$

and

$$
0<1-(A / J)=: \alpha_{2}<1
$$

In the second part of the theorem we may want $J$ even larger.
Let $\left(X_{0},\|\cdot\|\right)$ be the Banach space of bounded continuous functions $\phi:[0, \infty) \rightarrow \Re$ such that $\phi(t) \rightarrow 0$ as $t \rightarrow \infty$. For fixed $x(0)$, define $P: X_{0} \rightarrow X_{0}$ by $\phi \in X_{0}$ implies that

$$
(P \phi)(t)=z(t)+F(t)+\int_{0}^{t} R(t-s)\left[x(s)-\frac{G(s, x(s))}{J}\right] d s
$$

and note that $(P \phi)(t) \rightarrow 0$ as $t \rightarrow \infty$. Moreover, if $\phi, \eta \in X_{0}$ then
$|(P \phi)(t)-(P \eta)(t)| \leq \int_{0}^{t} R(t-s)\left|\phi(s)-\frac{G(s, \phi(s))}{J}-\eta(s)+\frac{G(s, \eta(s))}{J}\right| d s$.
This is a contraction with constant $\alpha=\max \left[\alpha_{1}, \alpha_{2}\right]=\alpha_{2}$ so $P$ has a unique fixed point $\phi \in X_{0}$. That proves the first part of the theorem.

For the second part, $z(t)=0$. Consider the aforementioned unique solution $\phi$ and find $T>0$ so that $s \geq T$ implies that

$$
\left|\phi(s)-\frac{G(s, \phi(s))}{J}\right| \leq \beta|\phi(s)| .
$$

Then interchanging the order of integration below, we obtain

$$
\begin{aligned}
\int_{0}^{t}|\phi(s)| d s & \leq \int_{0}^{t}|F(s)| d s+\int_{0}^{t} \int_{s}^{t} R(u-s) d u \mid \phi(s)-G(s, \phi(s) \mid d s \\
& \leq \int_{0}^{t}|F(s)| d s+\int_{0}^{T}|\phi(s)-G(s, \phi(s))| d s+\int_{T}^{t} \beta|\phi(s)| d s
\end{aligned}
$$

Hence,

$$
(1-\beta) \int_{T}^{t}|\phi(s)| d s \leq \int_{0}^{t}|F(s)| d s+\int_{0}^{T}|\phi(s)-G(s, \phi(s))| d s
$$

completing the proof.
There is a simple result which the reader should have in mind when viewing the nonlinear problems.

Proposition 3.6. Let $G(t) \geq 0$. Then

$$
L(t):=\int_{0}^{t} R(t-s) G(s) d s \in L^{1}[0, \infty) \Longleftrightarrow G \in L^{1}[0, \infty)
$$

Proof. If $G \in L^{1}$, clearly $L(t) \in L^{1}$. If $L(t) \in L^{1}$, then for large $t$

$$
\int_{0}^{t} \int_{0}^{u} R(u-s) G(s) d s d u \geq(1 / 2) \int_{0}^{t / 2} G(s) d s
$$

so $G \in L^{1}$.
The proposition reverses the classical theorem that the convolution of two $L^{1}$ functions is an $L^{1}$ function.

$$
\text { 4. } f \in L^{p}[0, \infty)
$$

This section is divided into two parts depending on the initial condition.
Case I: $x(0)=0$.
Return to (9)

$$
{ }^{c} D^{q} x=f(t)-a(t) x(t), \quad 0<q<1, \quad x(0)=0
$$

and (13)

$$
x(t)=F(t)-\int_{0}^{t} R(t-s)[\beta(s) x(s)] d s
$$

where

$$
F(t):=\frac{1}{J} \int_{0}^{t} R(t-s) f(s) d s
$$

and $R(t)$ satisfies (4)

$$
R(t)=C(t)-\int_{0}^{t} C(t-s) R(s) d s
$$

with $C$ defined in (3) as

$$
C(t)=\frac{1}{\Gamma(q)} J t^{q-1} \text { for } t>0
$$

We need the following lemma.
Lemma 4.1. If $f \in L^{p}$ then $F \in L^{p}$ for any $p \geq 1$ with $\|F\|_{p} \leq\|f\|_{p} / J$ and $F(t) \rightarrow 0$ as $t \rightarrow \infty$.

Proof. By Hölder's inequality, we have

$$
\begin{aligned}
\int_{0}^{t} R(t-s)|f(s)| d s & \leq\left(\int_{0}^{t} R(t-s) d s\right)^{\frac{p-1}{p}}\left(\int_{0}^{t} R(t-s)|f(s)|^{p} d s\right)^{\frac{1}{p}} \\
& \leq\left(\int_{0}^{t} R(t-s)|f(s)|^{p} d s\right)^{\frac{1}{p}}
\end{aligned}
$$

since $\int_{0}^{t} R(s) d s \leq 1$. This implies that

$$
\begin{aligned}
\int_{0}^{t}|F(u)|^{p} d u & \leq J^{-p} \int_{0}^{t} \int_{0}^{u} R(u-s)|f(s)|^{p} d s d u \\
& \leq J^{-p} \int_{0}^{t} \int_{s}^{t} R(u-s) d u|f(s)|^{p} d s \\
& =J^{-p} \int_{0}^{t} \int_{0}^{t-s} R(u) d u|f(s)|^{p} d s \\
& \leq J^{-p} \int_{0}^{\infty}|f(s)|^{p} d s<\infty
\end{aligned}
$$

and thus, $F \in L^{p}$ for any $p \geq 1$ and $\|F\|_{p} \leq\|f\|_{p} / J$. The fact that $F(t) \rightarrow 0$ as $t \rightarrow \infty$ follows from the inequality

$$
|F(t)|^{p} \leq J^{-p} \int_{0}^{t} R(t-s)|f(s)|^{p} d s
$$

the convolution of an $L^{1}$ function with a function tending to zero.

Corollary 4.2. For any $p \geq 1$,

$$
\int_{0}^{t} R(t-s)|f(s)| d s \leq\left(\int_{0}^{t} R(t-s)|f(s)|^{p} d s\right)^{1 / p}
$$

Theorem 4.3. If (11) holds and if $f \in L^{p}[0, \infty)$ for $p \geq 1$, then the solution of (9) is in $L^{p}[0, \infty)$ with

$$
\|x\|_{p} \leq\|f\|_{p} /[J(1-\alpha)] .
$$

Proof. We integrate (13) on $[0, t]$ and use the triangle inequality of the $L^{p}$-norm (or Minkowski inequality) to obtain

$$
\begin{aligned}
\left(\int_{0}^{t}|x(u)|^{p} d u\right)^{1 / p} \leq & \left(\int_{0}^{t}|F(u)|^{p} d u\right)^{1 / p} \\
& +\left[\int_{0}^{t}\left(\int_{0}^{u} R(u-s)|\beta(s) x(s)| d s\right)^{p} d u\right]^{1 / p} \\
\leq & \left(\int_{0}^{t}|F(u)|^{p} d u\right)^{1 / p} \\
& +\alpha\left(\left.\int_{0}^{t} \int_{0}^{u} R(u-s)| | x(s)\right|^{p} d s d u\right)^{1 / p}
\end{aligned}
$$

Interchange the order of integration and use (6), the property for $R(t)$, leaving us with

$$
(1-\alpha)\left(\int_{0}^{t}|x(s)|^{p} d s\right)^{1 / p} \leq\left(\int_{0}^{t}|F(s)|^{p} d s\right)^{1 / p}
$$

or

$$
\|x\|_{p} \leq\|f\|_{p} /[J(1-\alpha)] .
$$

This completes the proof.

Case II: $x(0) \neq 0$.
Let $\widetilde{R}(t)=1-\int_{0}^{t} R(s) d s$ and $z(t)=x(0) \widetilde{R}(t)$. For $x(0) \neq 0$, we write the counterpart of (9) as

$$
\begin{equation*}
{ }^{c} D^{q} x=f(t)-a(t) x(t), \quad 0<q<1, \quad x(0)=x_{0} . \tag{*}
\end{equation*}
$$

With $\beta(s)$ defined in (13), this can be inverted as

$$
x(t)=x(0)-\frac{1}{\Gamma(q)} \int_{0}^{t} J(t-s)^{q-1}\left[x(s)+\beta(s) x(s)-\frac{f(s)}{J}\right] d s .
$$

Then $x(t)$ solves

$$
\begin{equation*}
x(t)=z(t)+F(t)-\int_{0}^{t} R(t-s)[\beta(s) x(s)] d s \tag{14}
\end{equation*}
$$

Lemma 4.4. Let $R(t)$ be defined in (4). If $\widetilde{R}(t)=1-\int_{0}^{t} R(u) d u$, then

$$
\begin{equation*}
\widetilde{R}(t)=1-\int_{0}^{t} C(t-s) \widetilde{R}(s) d s \tag{15}
\end{equation*}
$$

Proof. We integrate (4) from 0 to $t$ to obtain

$$
\begin{aligned}
\int_{0}^{t} R(u) d u & =\int_{0}^{t} C(u) d u-\int_{0}^{t} \int_{0}^{u} C(u-s) R(s) d s d u \\
& =\int_{0}^{t} C(u) d u-\int_{0}^{t} \int_{0}^{t-s} C(u) d u R(s) d s
\end{aligned}
$$

This is equivalent to

$$
1-\int_{0}^{t} R(u) d u=1-\int_{0}^{t} C(u) d u+\int_{0}^{t} \int_{0}^{t-s} C(u) d u R(s) d s
$$

Taking into account that $\widetilde{R}^{\prime}(t)=-R(t)$, we arrive at
$\widetilde{R}(t)=1-\int_{0}^{t} C(u) d u-\int_{0}^{t} \int_{0}^{t-s} C(u) d u \widetilde{R}^{\prime}(s) d s$
(integrate the last term by parts)

$$
\begin{aligned}
& =1-\int_{0}^{t} C(u) d u-\left[\left.\int_{0}^{t-s} C(u) d u \widetilde{R}(s)\right|_{s=0} ^{s=t}+\int_{0}^{t} C(t-s) \widetilde{R}(s) d s\right] \\
& =1-\int_{0}^{t} C(u) d u-\left[-\int_{0}^{t} C(u) d u+\int_{0}^{t} C(t-s) \widetilde{R}(s) d s\right]
\end{aligned}
$$

as required.
Lemma 4.5. Let $R(t)$ be defined in (4) and $\widetilde{R}$ be defined in (15). Then $\widetilde{R}(t) \in L^{k}[0, \infty)$ for $k \geq 1$ if and only if $k>1 / q$.

Proof. We take Laplace transform of (15) to obtain

$$
\mathcal{L}(\widetilde{R})=\mathcal{L}(1)-\mathcal{L}(C) \mathcal{L}(\widetilde{R})
$$

and so

$$
\mathcal{L}(\widetilde{R})(s)=\frac{1}{s[1+\mathcal{L}(C)(s)]}
$$

We recall that $\mathcal{L}\left(t^{\nu}\right)(s)=\Gamma(\nu+1) s^{-\nu-1}$ for $\nu>-1$ and $\Re e(s)>0$ (see [7, p. 28]). Thus,

$$
\mathcal{L}(C)(s)=\frac{J}{\Gamma(q)} \mathcal{L}\left(t^{q-1}\right)(s)=\frac{J}{\Gamma(q)} \Gamma(q) s^{-q}=J s^{-q} .
$$

Substitute this equation into the right-hand side of $\mathcal{L}(\widetilde{R})(s)$ above to obtain

$$
\mathcal{L}(\widetilde{R})(s)=\frac{1}{s\left(1+J s^{-q}\right)}=\frac{s^{q-1}}{s^{q}+J}
$$

and hence,

$$
\widetilde{R}(t)=\mathcal{L}^{-1}\left[\frac{s^{q-1}}{s^{q}+J}\right]=E_{q, 1}\left(-J t^{q}\right)
$$

(see [7, p. 21]), where $E_{q, 1}$ is a member of the two parameter family of Mittag-Leffler functions (generalized exponential functions) defined as

$$
E_{\alpha, \beta}(z)=\sum_{k=0}^{\infty} \frac{z^{k}}{\Gamma(\alpha k+\beta)}, \alpha, \beta>0
$$

with the property, in particular, that

$$
\left|E_{q, 1}\left(-J t^{q}\right)\right| \leq \frac{\mu}{1+J t^{q}}
$$

for some constant $\mu$ and all $t \geq 0$ (see [7, p. 35]). Let $k>1 / q$ be fixed. Then

$$
|\widetilde{R}(t)|^{k} \leq(\mu / J)^{k} / t^{k q} \text { for } t \geq 1
$$

and thus, $\widetilde{R}(t) \in L^{k}[0, \infty)$ for all $k>1 / q$.
Conversely, suppose that $1 \leq k \leq 1 / q$. We show that $\widetilde{R}(t) \notin$ $L^{k}[0, \infty)$. In fact, it follows from an asymptotic expansion formula for $E_{\alpha, \beta}(z)$ (see [7, p. 33-34]) that

$$
E_{\alpha, \beta}(z)=-\sum_{k=1}^{p} \frac{z^{-k}}{\Gamma(\beta-\alpha k)}+O\left(|z|^{-1-p}\right) \quad \text { as } \quad|z| \rightarrow \infty
$$

where $0<\alpha<2, \beta>0, p \geq 1$ is an arbitrary integer, and $\mu \leq$ $|\arg (z)| \leq \pi$ with $\pi \alpha / 2<\mu<\min \{\pi, \pi \alpha\}$. Now for $p=1$, we have

$$
E_{q, 1}\left(-J t^{q}\right)=-\frac{1}{\Gamma(1-q)}\left(-J t^{q}\right)^{-1}+O\left(t^{-2 q}\right) \quad \text { as } t \rightarrow \infty
$$

Thus, there exists $\eta>0$ and $T>0$ such that

$$
\widetilde{R}(t)=E_{q, 1}\left(-J t^{q}\right) \geq \eta t^{-q} \quad \text { for } t \geq T
$$

and hence,

$$
|\widetilde{R}(t)|^{k} \geq \eta^{k} t^{-k q} \quad \text { for } t \geq T
$$

Since $0<k q \leq 1$, we see that $\widetilde{R} \notin L^{k}[0, \infty)$. This completes the proof.

Theorem 4.6. Suppose that (11) holds and $f \in L^{p}[0, \infty)$ for $p \geq 1$. Then the solution $x(t)$ of $\left(9^{*}\right)$ with $x(0) \neq 0$ is in $L^{p}[0, \infty)$ if and only if $p>1 / q$. Moreover, if $p>1 / q$, then

$$
\begin{equation*}
\|x\|_{p} \leq L\left[|x(0)|\|\widetilde{R}\|_{p}+\|f\|_{p}\right] \tag{16}
\end{equation*}
$$

for some constant $L>0$.
Proof. If $p>1 / q$, we have $\widetilde{R} \in L^{p}[0, \infty)$ by Lemma 4.5. It then follows from (14) that

$$
\begin{aligned}
\|x\|_{p} & \leq\|z\|_{p}+\|F\|_{p}+\alpha\|x\|_{p} \\
& \leq|x(0)|\|\widetilde{R}\|_{p}+\|f\|_{p} / J+\alpha\|x\|_{p}
\end{aligned}
$$

This implies that $x \in L^{p}[0, \infty)$ and (16) holds.
Conversely, suppose that $1 \leq p \leq 1 / q$. Then $\widetilde{R} \notin L^{p}[0, \infty)$ by Lemma 4.5. We write (14) as

$$
\widetilde{R}(t) x(0)=x(t)-F(t)+\int_{0}^{t} R(t-s)[\beta(s) x(s)] d s
$$

If $x \in L^{p}[0, \infty)$, then

$$
\|\widetilde{R}\|_{p} \leq\left[\|x\|_{p}+\|f\|_{p} / J+\alpha\|x\|_{p}\right] /|x(0)|<\infty
$$

which yields $\widetilde{R} \in L^{p}[0, \infty)$, a contradiction.

Remark: If $f \equiv 0$, then by (16), the zero solution of $\left(9^{*}\right)$ is $L^{p_{-}}$ asymptotically stable. We also observe from (16) that solutions of $\left(9^{*}\right)$ are uniformly bounded and uniformly ultimately bounded in $L^{p}[0, \infty)$ for $p>1 / q$ at $t=0$.

We consider the nonlinear equation

$$
\begin{equation*}
{ }^{c} D^{q} x=-a(t) x^{3}(t)+f(t), \quad 0<q<1 \tag{17}
\end{equation*}
$$

where $a, f:[0, \infty) \rightarrow \Re$ are continuous and there are positive numbers $\epsilon$ and $M$ such that (11) holds. We then find $J>0, \alpha<1$ with $|J-a(t)| \leq \alpha J$. We invert (17) as

$$
x(t)=x(0)-\frac{1}{\Gamma(q)} \int_{0}^{t}(t-s)^{q-1}\left[a(s) x^{3}(s)-f(s)\right] d s
$$

which we write

$$
x(t)=x(0)-\frac{1}{\Gamma(q)} \int_{0}^{t}(t-s)^{q-1}\left[J x(t)-J x(t)+a(s) x^{3}(s)-f(s)\right] d s
$$

We now decompose it into

$$
z(t)=x(0)-\frac{J}{\Gamma(q)} \int_{0}^{t}(t-s)^{q-1} z(s) d s
$$

with solution $z(t)=x(0) \widetilde{R}(t)$ and
(18) $x(t)=z(t)+\int_{0}^{t} R(t-s)\left[x(s)-x^{3}(s)+\frac{(J-a(s))}{J} x^{3}(s)\right] d s+F(t)$
where $R(t)$ is defined in (4) and

$$
F(t):=\frac{1}{J} \int_{0}^{t} R(t-s) f(s) d s
$$

Note that $F \in L^{p}$ if $f \in L^{p}$ for $p \geq 1$ and $F(t) \rightarrow 0$ as $t \rightarrow 0$.
Theorem 4.7. Suppose (11) holds. If $f \in L^{p}[0, \infty)$ for $p \geq 1$ and $|F(t)|<(1-\alpha) \eta^{3}$ with $0<\eta<\sqrt{3} / 3$ and $\left(1-\eta^{2}\right)^{3(p-1)}>\alpha$, then the solution of (17) with $x(0)=0$ is in $L^{3 p}[0, \infty)$ with $\|x\|_{3 p} \leq L\|f\|_{p}^{1 / 3}$ for some constant $L>0$.

Proof. For $x(0)=0$, we see that $z(t)=0$ and (18) becomes

$$
\begin{equation*}
x(t)=\int_{0}^{t} R(t-s)\left[x(s)-x^{3}(s)+\frac{(J-a(s))}{J} x^{3}(s)\right] d s+F(t) \tag{19}
\end{equation*}
$$

First, we claim that $|x(t)|<\eta$ for all $t \geq 0$. If there is a $\bar{t}>0$ such that $|x(\bar{t})|=\eta$ and $|x(s)|<|x(\bar{t})|$ for all $0 \leq s<\bar{t}$, then

$$
\begin{aligned}
|x(\bar{t})| & \leq \int_{0}^{\bar{t}} R(\bar{t}-s)\left[\left(|x(s)|-|x(s)|^{3}\right)+\alpha|x(s)|^{3}\right] d s+|F(\bar{t})| \\
& \leq\left[|x(\bar{t})|-|x(\bar{t})|^{3}+\alpha|x(\bar{t})|^{3}\right] \int_{0}^{\bar{t}} R(\bar{t}-s) d s+|F(\bar{t})|
\end{aligned}
$$

Since $\eta-(1-\alpha) \eta^{3}>0$, this yields

$$
(1-\alpha) \eta^{3} \leq|F(\bar{t})|
$$

which contradicts the condition on $F$. Thus, $|x(t)|<\eta$ for all $t \geq 0$.
Next, we set $\gamma=3(p-1)$ and multiply (19) by $\left(|x(t)|-|x(t)|^{3}\right)^{\gamma}$ to obtain

$$
\begin{align*}
& \left(|x(t)|-|x(t)|^{3}\right)^{\gamma+1}+|x(t)|^{3}\left(|x(t)|-|x(t)|^{3}\right)^{\gamma}  \tag{20}\\
& \leq \int_{0}^{t} R(t-s)\left[|x(s)|-|x(s)|^{3}\right] d s\left(|x(t)|-|x(t)|^{3}\right)^{\gamma} \\
& +\alpha \int_{0}^{t} R(t-s)|x(s)|^{3} d s\left(|x(t)|-|x(t)|^{3}\right)^{\gamma} \\
& +|F(t)|\left(|x(t)|-|x(t)|^{3}\right)^{\gamma}
\end{align*}
$$

Use Young's inequality to write

$$
\begin{aligned}
& \left(|x(s)|-|x(s)|^{3}\right)\left(|x(t)|-|x(t)|^{3}\right)^{\gamma} \\
& \leq \frac{1}{\gamma+1}\left(|x(s)|-|x(s)|^{3}\right)^{\gamma+1}+\frac{\gamma}{\gamma+1}\left(|x(t)|-|x(t)|^{3}\right)^{\gamma+1}
\end{aligned}
$$

and

$$
\begin{aligned}
& |x(s)|^{3}\left(|x(t)|-|x(t)|^{3}\right)^{\gamma} \\
& \quad \leq \frac{3}{\gamma+3}|x(s)|^{\gamma+3}+\frac{\gamma}{\gamma+3}\left(|x(t)|-|x(t)|^{3}\right)^{\gamma+3}
\end{aligned}
$$

so that

$$
|F(t)|\left(|x(t)|-|x(t)|^{3}\right)^{\gamma} \leq K|F(t)|^{p}+\delta\left(|x(t)|-|x(t)|^{3}\right)^{\gamma+3}
$$

where $0<\delta<\left(1-\eta^{2}\right)^{\gamma}-\alpha$ and $K=K(\delta)$ is a constant.
Substitute these inequalities into (20) and integrate on $[0, t]$ to obtain

$$
\begin{aligned}
& \int_{0}^{t}\left(|x(u)|-|x(u)|^{3}\right)^{\gamma+1} d u+\int_{0}^{t}|x(u)|^{3}\left(|x(u)|-|x(u)|^{3}\right)^{\gamma} d u \\
& \leq \frac{1}{\gamma+1} \int_{0}^{t} \int_{0}^{u} R(u-s)\left(|x(s)|-|x(s)|^{3}\right)^{\gamma+1} d s d u \\
& +\frac{\gamma}{\gamma+1} \int_{0}^{t} \int_{0}^{u} R(u-s)\left(|x(u)|-|x(u)|^{3}\right)^{\gamma+1} d s d u \\
& +\frac{3 \alpha}{\gamma+3} \int_{0}^{t} \int_{0}^{u} R(u-s)|x(s)|^{\gamma+3} d s d u \\
& +\frac{\gamma \alpha}{\gamma+3} \int_{0}^{t} \int_{0}^{u} R(u-s)\left(|x(u)|-|x(u)|^{3}\right)^{\gamma+3} d s d u \\
& +K \int_{0}^{t}|F(u)|^{p} d u+\delta \int_{0}^{t}\left(|x(u)|-|x(u)|^{3}\right)^{\gamma+3} d u
\end{aligned}
$$

Interchange the order of integration, use the property $\int_{0}^{t} R(u) d u \leq 1$, and cancel terms, leaving us with

$$
\begin{align*}
& \int_{0}^{t}|x(u)|^{3}\left(|x(u)|-|x(u)|^{3}\right)^{\gamma} d u  \tag{21}\\
& \leq \frac{3 \alpha}{\gamma+3} \int_{0}^{t}|x(s)|^{\gamma+3} d s+\frac{\gamma \alpha}{\gamma+3} \int_{0}^{t}\left(|x(u)|-|x(u)|^{3}\right)^{\gamma+3} d u \\
& +K \int_{0}^{t}|F(u)|^{p} d u+\delta \int_{0}^{t}\left(|x(u)|-|x(u)|^{3}\right)^{\gamma+3} d u \\
& \leq(\alpha+\delta) \int_{0}^{t}|x(s)|^{\gamma+3} d s+K \int_{0}^{t}|F(u)|^{p} d u
\end{align*}
$$

The last line was obtained by noting that $|x(t)|>|x(t)|^{3}$. Since $|x(t)|<$ $\eta$ for all $t \geq 0$, we have

$$
\begin{aligned}
& \left(1-\eta^{2}\right)^{\gamma} \int_{0}^{t}|x(s)|^{\gamma+3} d s \\
& \leq \int_{0}^{t}|x(u)|^{\gamma+3}\left(1-|x(u)|^{2}\right)^{\gamma} d u \\
& =\int_{0}^{t}|x(u)|^{3}\left(|x(u)|-|x(u)|^{3}\right)^{\gamma} d u \\
& \leq(\alpha+\delta) \int_{0}^{t}|x(s)|^{\gamma+3} d s+K \int_{0}^{t}|F(u)|^{p} d u
\end{aligned}
$$

Note that $\gamma+3=3 p$ and $\left(1-\eta^{2}\right)^{\gamma}>\alpha+\delta$. Moving the next-to-last integral to the left-hand-side of the display, we see that $x \in L^{3 p}[0, \infty)$ and $\|x\|_{3 p} \leq L\|f\|_{p}^{1 / 3}$ for some constant $L>0$. This completes the proof.

Corollary 4.8. Suppose (11) holds. If $f \in L^{p}[0, \infty)$ for $p \geq 1$, and if $f$ is bounded on $[0, \infty)$, then the solution of (17) with $x(0)=0$ is in $L^{3 p}[0, \infty)$ with $\|x\|_{3 p} \leq L\|f\|_{p}^{1 / 3}$ for some constant $L>0$.
Proof. Since $f$ is bounded on $[0, \infty)$, there exists a constant $H>0$ such that $|f(t)| \leq H$ for all $t \geq 0$. For the fixed $0<\alpha<1$ and $p \geq 1$, we choose $\eta>0$ so that $0<\eta<\sqrt{3} / 3$ and $\left(1-\eta^{2}\right)^{3(p-1)}>\alpha$. We see from the proof of Lemma 4.1 that

$$
\begin{aligned}
|F(t)|^{p} & \leq J^{-p} \int_{0}^{t} R(t-s)|f(s)|^{p} d s \\
& \leq J^{-p} H^{p} \int_{0}^{t} R(t-s) d s \leq(H / J)^{p} \leq(2 H / M)^{p}
\end{aligned}
$$

by the definition of $J$, where $M$ is defined in (11). This yields

$$
|F(t)| \leq 2 H / M
$$

We may choose $M$ sufficiently large so that $2 H / M<(1-\alpha) \eta^{3}$, and thus

$$
|F(t)| \leq(1-\alpha) \eta^{3} .
$$

Therefore, all conditions of Theorem 4.7 are satisfied, and the assertion of Corollary 4.8 follows.

We now consider the case $x(0) \neq 0$.

Theorem 4.9. Suppose (11) holds. If $f \in L^{p}[0, \infty)$ for $p>1 / q$ and $|F(t)|<(1-\alpha) \eta^{3}$ with $0<\eta<\sqrt{3} / 3$ and $\left(1-\eta^{2}\right)^{3(p-1)}>\alpha$, then the solution of (17) is in $L^{3 p}[0, \infty)$ with

$$
\begin{equation*}
\|x\|_{3 p} \leq L\left[|x(0)|\|\widetilde{R}\|_{p}+\|f\|_{p}\right]^{1 / 3} \tag{22}
\end{equation*}
$$

for some constant $L>0$.

Proof. The proof is similar to that of Theorem 4.7 (with an additional term) by working with (18). Since $F(t) \rightarrow 0$ as $t \rightarrow \infty$, we see that $\|F\|<(1-\alpha) \eta^{3}$ and $|x(t)|<\eta$ for all $t \geq 0$ if

$$
|x(0)|+\|F\|<(1-\alpha) \eta^{3} .
$$

We now set $\gamma=3(p-1)$, multiply (18) by $\left(|x(t)|-|x(t)|^{3}\right)^{\gamma}$, and follow through the calculations in (20) and (21) to obtain

$$
\begin{aligned}
& \int_{0}^{t}|x(u)|^{3}\left(|x(u)|-|x(u)|^{3}\right)^{\gamma} d u \\
& \leq(\alpha+\delta) \int_{0}^{t}|x(s)|^{\gamma+3} d s+K \int_{0}^{t}|F(u)|^{p} d u \\
& \quad+|x(0)| \int_{0}^{t} \widetilde{R}(u)\left(|x(u)|-|x(u)|^{3}\right)^{\gamma} d u
\end{aligned}
$$

Next, find $0<\delta_{1}<\left(1-\eta^{2}\right)^{\gamma}-(\alpha+\delta)$ and $K_{1}>0$ such that

$$
\begin{aligned}
& |x(0)| \int_{0}^{t} \widetilde{R}(u)\left(|x(u)|-|x(u)|^{3}\right)^{\gamma} d u \\
& \leq|x(0)|^{p} K_{1} \int_{0}^{t}|\widetilde{R}(u)|^{p} d u+\delta_{1} \int_{0}^{t}\left(|x(u)|-|x(u)|^{3}\right)^{\gamma+3} d u .
\end{aligned}
$$

Since $|x(t)|<\eta$ for all $t \geq 0$, we have

$$
\begin{aligned}
& \left(1-\eta^{2}\right)^{\gamma} \int_{0}^{t}|x(s)|^{\gamma+3} d s \\
& \leq \int_{0}^{t}|x(u)|^{\gamma+3}\left(1-|x(u)|^{2}\right)^{\gamma} d u \\
& =\int_{0}^{t}|x(u)|^{3}\left(|x(u)|-|x(u)|^{3}\right)^{\gamma} d u \\
& \leq \\
& \left(\alpha+\delta+\delta_{1}\right) \int_{0}^{t}|x(s)|^{\gamma+3} d s \\
& \quad+|x(0)|^{p} K_{1} \int_{0}^{t}|\widetilde{R}(u)|^{p} d u+K \int_{0}^{t}|F(u)|^{p} d u .
\end{aligned}
$$

This implies that there exists a constant $K_{2}>0$ such that

$$
\begin{equation*}
\int_{0}^{t}|x(s)|^{\gamma+3} d s \leq K_{2}\left[|x(0)|^{p} \int_{0}^{t}|\widetilde{R}(u)|^{p} d u+\int_{0}^{t}|F(u)|^{p} d u\right] \tag{23}
\end{equation*}
$$

for all $t \geq 0$. Raising both sides of (23) to the power $1 /(\gamma+3)$ and taking into account that $\gamma+3=3 p$ and $\widetilde{R} \in L^{p}$, we see

$$
\begin{aligned}
\|x\|_{3 p} & \leq K_{2}^{\frac{1}{3 p}}\left[|x(0)|^{p}\|\widetilde{R}\|_{p}^{p}+\|F\|_{p}^{p}\right]^{\frac{1}{3 p}} \\
& \leq K_{2}^{\frac{1}{3 p}} 2^{\frac{1}{3 p}}\left[|x(0)|\|\widetilde{R}\|_{p}+\|f\|_{p} / J\right]^{\frac{1}{3}} \\
& \leq K_{2}^{\frac{1}{3 p}} 2^{\frac{1}{3 p}}(1+1 / J)^{\frac{1}{3}}\left[|x(0)|\|\widetilde{R}\|_{p}+\|f\|_{p}\right]^{\frac{1}{3}} \\
& =: L\left[|x(0)|\|\widetilde{R}\|_{p}+\|f\|_{p}\right]^{\frac{1}{3}}
\end{aligned}
$$

This completes the proof.

Corollary 4.10. Suppose (11) holds. If $f \in L^{p}[0, \infty)$ for $p>1 / q$, and if $f$ is bounded on $[0, \infty)$, then the solution of (17) is in $L^{3 p}[0, \infty)$, and it satisfies (22).

Remark: If $f \equiv 0$, then by (22), the zero solution of (17) is $L^{3 p_{-}}$ asymptotically stable. One may also note that Theorem 4.7 and Theorem 4.9 are general results. Everything would work for

$$
{ }^{c} D^{q} x=-a(t) x^{2 n+1}(t)+f(t), \quad 0<q<1
$$

for a positive integer $n$.

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## References

[1] L. C. Becker, Resolvents and solutions of weakly singular linear Volterra integral equations, Nonlinear Anal.74(2011), 1892-1912.
[2] T. A. Burton, Stability by Fixed Point Theory for Functional Differential Equations, Dover, Mineola, New York, 2006.
[3] Kai Diethelm, The Analysis of Fractional Differential Equations, Springer, Heidelberg, 2010.
[4] V. Lakshmikantham, S. Leela, J. Vasundhara Devi, Theory of Fractional Dynamic Systems, Cambridge Scientific Publishers, Cottenham, Cambridge, 2009.
[5] V. Lakshmikantham and A. S. Vatsala, Basic theory of fractional differential equations, Nonlinear Anal. 69(2008), 2677-2682.
[6] R. K. Miller, Nonlinear Integral Equations, Benjamin, Menlo Park, CA, 1971.
[7] I. Podlubny, Fractional Differential Equations, Academic Press, San Diego, 1999.
[8] A. Windsor, A contraction mapping proof of the smooth dependence on parameters of solutions to volterra integral equations, Nonlinear Anal. 72(2010), 3627-3634.

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